

- Parity  $P$  acts as  $P = P^\dagger = P^{-1}$  with  $P\vec{v}P = -\vec{v}$  and  $P\vec{a}P = \vec{a}$  for vectors (e.g.  $\vec{x}$  and  $\vec{p}$ ) and axial pseudo-vectors (e.g.  $\vec{L}$  and  $\vec{B}$ ) respectively. Note that it is not a rotation. Parity preserves all the basic commutation relations, but it is not necessarily a symmetry of the Hamiltonian. The names scalar and pseudo scalar are often used to denote parity even vs odd:  $POP = \mathcal{O}$  for “scalar” operators (like  $\vec{x}^2$  or  $\vec{L}^2$ , and  $PO'P = -\mathcal{O}'$  for pseudo scalar operators (like  $\vec{S} \cdot \vec{x}$ ).

- Nature respects parity aside from some tiny effects at the fundamental level. E.g. in electricity and magnetism, we can preserve Maxwell’s equations via  $\vec{V} \rightarrow \vec{V}$  and  $\vec{V}_{axial} \rightarrow +\vec{V}_{axial}$ .

- In QM, parity is represented by a unitary operator such that  $P|\vec{r}\rangle = |-\vec{r}\rangle$  and  $P^2 = 1$ . So the eigenvectors of  $P$  have eigenvalue  $\pm 1$ . A state is even or odd under parity if  $P|\psi\rangle = \pm|\psi\rangle$ ; in position space,  $\psi(-\vec{x}) = \pm\psi(\vec{x})$ . In radial coordinates, parity takes  $\cos\theta \rightarrow -\cos\theta$  and  $\phi \rightarrow \phi + \pi$ , so  $P|\alpha, \ell, m\rangle = (-1)^\ell|\alpha, \ell, m\rangle$ . This is clear also from the  $\ell = 1$  case, any getting general  $\ell$  from tensor products.

If  $[H, P] = 0$ , then the non-degenerate energy eigenstates can be written as  $P$  eigenstates.

Parity in 1d:  $PxP = -x$  and  $PpP = -p$ . The Hamiltonian respects parity if  $V(-x) = V(x)$ . The energy eigenstates can thus be chosen to be even or odd. E.g. 1d particle in a box. E.g. for 1d, SHO the state  $|n\rangle$  has parity  $(-1)^n$ , since the  $a$  and  $a^\dagger$  operators are parity odd. Similarly for the 3d SHO: the state with quantum numbers  $|n_x, n_y, n_z\rangle$  has parity  $(-1)^{n_x+n_y+n_z}$ . This agrees with the fact that the parity is  $(-1)^\ell$ , since  $\ell = n \bmod 2$ .

Recall example of bound states in a potential well,  $V(x) = -V_0\theta(\frac{1}{2}a - |x|)$ . Write down even and odd parity solutions, and the conditions from having  $\psi$  and  $\psi'$  continuous. Even parity solutions thus have  $p \tan(pa/2\hbar) = \hbar\kappa$  with  $p \equiv \sqrt{2m(E + |V|)}$  and  $\hbar\kappa = \sqrt{-2mE}$ , where  $E < 0$ . Odd parity solutions have  $p \cot(pa/2\hbar) = -\hbar\kappa$ . For a shallow  $V_0$ , only one even solution, then get alternating even / odd solutions as  $V_0$  increases.

E.g. symmetric double well, or SHO-type double well:  $E_A > E_S$ . Double well example (mention  $e^{-S_{Eucl}/\hbar}$  tunneling and instantons).

- Parity selection rules: if  $[H, P] = 0$ , time evolution preserves parity and the states and operators can be assigned definite parity. Show that  $\langle \chi | \mathcal{O} | \psi \rangle = 0$  unless  $(-1)^{P_\chi + P_\mathcal{O} + P_\psi} = 1$ , where  $(-1)^{P_{\chi, \psi, \mathcal{O}}}$  are the parity signs of the operators and states. Illustrate with SHO examples.