

4/10/18 Lecture outline

★ Reading: Zwiebach chapters 1 and 2.

• Last time: Statement of relativity principle: physics is indistinguishable among all inertial frames. If one frame is inertial, all frames moving relative to that at a constant velocity are also inertial. Related by  $x^{\mu'} = \Lambda^{\mu'}_{\nu} x^{\nu}$  where  $\Lambda^T \eta \Lambda = \eta$ , where  $\eta$  is the metric, chosen to have mostly plus convention. All 4-vectors transform the same way:  $b^{\mu'} = \Lambda^{\mu'}_{\nu} b^{\nu}$ . A tensor transforms as  $T^{\mu'\nu'} = \Lambda^{\mu'}_{\kappa} \Lambda^{\nu'}_{\sigma} T^{\kappa\sigma}$ . Can contract or raise or lower indices with  $\eta_{\mu\nu}$ , e.g.  $a^{\mu} b^{\nu} \eta_{\mu\nu} = -a^0 b^0 + \vec{a} \cdot \vec{b}$  and  $T^{\mu}_{\mu} = -T^{00} + \delta_{ij} T^{ij}$  are 4-scalars. 4-scalars are invariant under Lorentz transformation. The proper time element  $d\tau = \sqrt{-dx^{\mu} dx_{\mu}/c^2}$  is a 4-scalar.  $u^{\mu} = dx^{\mu}/d\tau$  is the velocity 4-vector, and  $d^2 x^{\mu}/d\tau^2$  is the acceleration 4-vector. The statement of relativity requires that all formulas relate quantities that transform the same way, for example  $\vec{F} = m\vec{a}$  can fit with relativity if it becomes a 4-vector equation:  $f^{\mu} = mdp^{\mu}/d\tau$  (the time component here relates power to the time derivative of energy).

$p^{\mu} = (E/c, p_x, p_y, p_z)$ , with  $p_{\mu} p^{\mu} = -m^2 c^2$ .  $p^{\mu}$  transforms as a Lorentz 4-vector,  $p^{\mu'} = \Lambda^{\mu'}_{\nu} p^{\nu}$ . Proper time:  $ds^2 = c^2 dt_p^2 = c^2 dt^2 (1 - \beta^2)$ .  $u^{\mu} = cdx^{\mu}/ds = dx^{\mu}/dt_p = \gamma(c, \vec{v})$ , and  $u_{\mu} u^{\mu} = -c^2$ . A massive point particle has  $p^{\mu} = mu^{\mu}$ . Massless particles, like the photon, have  $p^{\mu}$  with  $p^{\mu} p_{\mu} = 0$ .

• Quantum mechanics: replace  $p^{\mu} = (H/c, \vec{p}) \rightarrow -i\hbar\partial^{\mu}$ . Free particle wavefunction  $\psi \sim \exp(ip \cdot x/\hbar)$ ;  $p_{\mu} x^{\mu} \equiv p \cdot x$  is Lorentz invariant.

• Light cone coordinates:  $x^{\pm} = \frac{1}{\sqrt{2}}(x^0 \pm x^1)$ . The bad: spoils rotational symmetry. The good: will make it easier to quantize string theory (avoids having to discuss here more advanced alternatives, which do not require going to the light cone coordinates).  $-ds^2 = -2dx^+ dx^- + dx_2^2 + dx_3^2 = -\hat{\eta}_{\mu\nu} dx^{\mu} dx^{\nu}$ .  $a_{\pm} = -a^{\mp}$ . Take  $p^{\pm} = \frac{1}{\sqrt{2}}(p^0 \pm p^1) = -p_{\mp}$ . So  $i\hbar\partial_{x^+} \rightarrow -p_+ = E_{lc}/c$ , i.e.  $p^- = E_{lc}/c$ .

• Extra (spacelike) dimensions, e.g. 2 extra dimensions:  $-ds^2 = -c^2 dt^2 + \sum_{i=1}^5 (dx^i)^2$ . Consider one extra space dimension, taken to be a circle,  $x \sim x + 2\pi R$ . Now consider  $(x, y) \sim (x + 2\pi R, y) \sim (x, y + 2\pi R)$ ; gives a torus. Orbifold, e.g.  $z \sim e^{i\pi i/N} z$ , gives a cone (singular at fixed point).

• Recall QM:  $[x^i, p_j] = i\hbar\delta_j^i$ . Particle in square well box of size  $a$ :  $E = (n\pi/a)^2/2m$ . Now particle in periodic box,  $x_4 \sim x_4 + 2\pi R$ . The other directions,  $x^{\mu}$ , are given by some standard Hamiltonian, e.g. the hydrogen atom, which we'll call  $H_{4d}$ . So  $H_{5d} = H_{4d} + \hat{p}_4^2/2m$ , with  $\hat{p}_4 = -i\hbar\partial_{x_4}$  in position space. The 4d energy eigenstates are then given by separation of variables to be  $\psi_{E_{5d}}(\vec{x}, x_4) = \psi_{E_{4d}}(\vec{x}) \frac{1}{\sqrt{2\pi R}} e^{i\ell x_4/R}$ , with  $\ell$  an integer, and  $\psi_{E_{4d}}$  is an energy eigenstate of the 4d problem. So  $E_{5d} = E_{4d} + \ell^2/2mR^2$ . For  $R$  small, the low energy states are simply those with  $\ell = 0$ , and the extra dimension is unseen.