

## 5/17/07 Lecture 13 outline

- In quantum mechanics, we replace physical observables, like  $x$ ,  $p$ ,  $E$ , etc. with Hermitian operators. The observed quantities are the eigenvalues. Write e.g.

$$\hat{x}|x\rangle = x|x\rangle, \quad \hat{p}|p\rangle = p|p\rangle, \quad H|E\rangle = E|E\rangle.$$

These operators generally act in an infinite dimensional space, the Hilbert space, but this generally doesn't complicate things much (from a physicist's perspective).

The  $x$  and  $p$  eigenvalues are continuous.  $E$  can be continuous (e.g. free particle), discrete (e.g. particle in a box with  $V = \infty$  outside, or more generally when  $V \rightarrow \infty$  for  $|x| \rightarrow \text{infy}$ ), or a combination of continuous and discrete (e.g. for the particle in a finite potential well).

- As we have discussed, the operators  $\hat{x}$  and  $\hat{p}$  satisfy  $[\hat{x}, \hat{p}] = i\hbar$ . The fact that they don't commute means that they don't have simultaneous eigenvectors, they can't be simultaneously diagonalized. Nevertheless, there is a close relation between them. For any ket  $\chi$ , we have

$$\langle x|\hat{p}|\chi\rangle = -i\hbar \frac{d}{dx} \langle x|\chi\rangle \quad \text{and} \quad \langle p|\hat{x}|\chi\rangle = i\hbar \frac{d}{dp} \langle p|\chi\rangle.$$

These relations are consistent with  $[\hat{x}, \hat{p}] = i\hbar$ .

Their separate eigenkets satisfy  $\langle x'|x\rangle = \delta(x - x')$ , and  $\langle p'|p\rangle = \delta(p - p')$ . The completeness relations are

$$\int_{-\infty}^{\infty} dx |x\rangle \langle x| = 1_{op} \quad \int_{-\infty}^{\infty} dp |p\rangle \langle p| = 1_{op}.$$

The relation between these bases is  $\langle x|p\rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{ipx/\hbar}$ . Note that this satisfies  $\langle x|p|p\rangle = (-i\hbar \frac{d}{dx}) \langle x|p\rangle = p \langle x|p\rangle$ .

The wavefunction of a QM system is represented by an abstract vector in the Hilbert space,  $|\psi\rangle$ . The wavefunction in position space is  $\psi(x) = \langle x|\psi\rangle$ . The wavefunction in momentum space is  $\phi(p) = \langle p|\psi\rangle$ . Explain the meaning of the Fourier transform between them, with  $\langle x|p\rangle$ .

- **Measurement (this is a key point!):** If measuring observable to operator  $A$ , write  $|\psi\rangle = \sum_i |a_i\rangle \langle a_i|\psi\rangle$  (using completeness). The probability to measure  $A = a_i$  in this state is then  $|\langle a_i|\psi\rangle|^2$ . Immediately after the measurement, the wavefunction collapses,  $|\psi\rangle \rightarrow |a_i\rangle$ . If operators  $A$  and  $B$  commute, they can be simultaneously diagonalized. Discuss measurement and operators which do, or do not, commute.

- The Schrodinger equation in this notation is

$$i\hbar \frac{\partial}{\partial t} |\psi(t)\rangle = H|\psi\rangle.$$

The way we wrote it before was in the  $x$  basis. Discuss it in other bases.

- E.g. for the particle in the infinite potential well, we have energy eigenstates  $|n\rangle$ ,  $n = 1, 2, \dots$ , with  $H|n\rangle = E_n|n\rangle$ , for  $E_n = n^2\pi^2\hbar^2/2mL^2$ . These energy eigenstates satisfy the usual relations for a complete, orthonormal basis:

$$\langle n|m\rangle = \delta_{nm} \quad \sum_{n=1}^{\infty} |n\rangle\langle n| = 1.$$

The relation between the  $|n\rangle$  and the  $|x\rangle$  basis is

$$\langle x|n\rangle = u_n(x) = \begin{cases} \sqrt{\frac{2}{L}} \sin(n\pi x/L) & \text{for } 0 \leq x \leq L \\ 0 & \text{otherwise} \end{cases}.$$