10/14 Lecture outline

- * Reading: Luke, chapter 5
- Last time, Dyson's formula: $S = Te^{-i\int d^4x \mathcal{H}_{int}}$. Use this to compute $\langle f|S-1|i\rangle$, using Wick's theorem,

$$T(\phi_1 \dots \phi_n) =: \phi_1 \dots \phi_n :+ \sum_{contractions} : \phi_1 \dots \phi_n :$$

to eliminate the T ordering

• Illustrate this for our simple example of interacting theory:

$$\mathcal{L} = \frac{1}{2}(\partial \phi^2 - \mu^2 \phi^2) + (\partial \psi^{\dagger} \partial \psi - m^2 \psi^{\dagger} \psi) - g \phi \psi \psi^{\dagger}.$$

Toy model for interacting nucleons and mesons. Treat last term as a perturbation.

Recall $\phi \sim a + a^{\dagger}$ for "mesons," $\psi \sim b + c^{\dagger}$, and $\psi^{\dagger} \sim b^{\dagger} + c$. We'll say that b annihilates a nucleon N and c^{\dagger} creates an anti-nucleon \bar{N} . Conservation law, conserved charge $Q = N_b - N_c$. Examples of states:

$$|\phi(p)\rangle = a^{\dagger}(p)|0\rangle, \qquad |N(p)\rangle = b^{\dagger}(p)|0\rangle, \qquad |\bar{N}(p)\rangle = c^{\dagger}(p)|0\rangle.$$

Note then e.g.

$$\langle 0|\phi(x)|\phi(p)\rangle = e^{-ip\cdot x}, \qquad \langle 0|\psi(x)|N(p)\rangle = e^{-ip\cdot x}, \qquad \langle 0|\psi^{\dagger}(x)|N(p)\rangle = 0.$$

Example: meson decay. $|i\rangle = a^{\dagger}(p)|0\rangle$, $|f\rangle = b^{\dagger}(q_1)c^{\dagger}(q_2)|0\rangle$. Compute $\langle f|S|i\rangle = -ig\delta^4(p-q_1-q_2)$ to $\mathcal{O}(g)$.

Now consider $N + N \to N + N$, to $\mathcal{O}(g^2)$. The initial and final states are

$$|i\rangle = b^{\dagger}(p_1)b^{\dagger}(p_2)|0\rangle, \qquad \langle f| = \langle 0|b(p_1')b(p_2').$$

The term that contributes to scattering at $\mathcal{O}(g^2)$ is

$$T\frac{(-ig)^2}{2!} \int d^4x_1 d^4x_2 \phi(x_1) \psi^{\dagger}(x_1) \psi(x_1) \phi(x_2) \psi^{\dagger}(x_2) \psi(x_2).$$

The term that contributes to S-1 thus involves

$$\langle p_1' p_2' | : \psi^{\dagger}(x_1) \psi(x_1) \psi^{\dagger}(x_2) \psi(x_2) : | p_1 p_2 \rangle = \langle p_1' p_2' | : \psi^{\dagger}(x_1) \psi^{\dagger}(x_2) | 0 \rangle \langle 0 | \psi(x_1) \psi(x_2) | p_1, p_2 \rangle.$$

$$= \left(e^{i(p_1'x_1+p_2'x_2)} + e^{i(p_1'x_2+p_2'x_1)}\right) \left(e^{-i(p_1x_1+p_2x_2)} + e^{-i(p_1x_2+p_2x_1)}\right).$$

The amplitude involves this times $D_F(x_1 - x_2)$ (from the contraction), with the prefactor and integrals as above. The final result is

$$i(-ig)^{2} \left[\frac{1}{(p_{1}-p'_{1})^{2}-\mu^{2}} + \frac{1}{(p_{1}-p'_{2})^{2}-\mu^{2}} \right] (2\pi)^{4} \delta^{(4)}(p_{1}+p_{2}-p'_{1}-p'_{2}).$$

Explicitly, in the CM frame, $p_1 = (\sqrt{p^2 + m^2}, e\hat{e})$ and $p_2 = (\sqrt{p^2 + m^2}, -p\hat{e})$, $p_1' = (\sqrt{p^2 + m^2}, p\hat{e}')$, $p_2' = (\sqrt{p^2 + m^2}, -p\hat{e}')$, where $\hat{e} \cdot \hat{e}' = \cos \theta$, and get

$$\mathcal{A} = g^2 \left(\frac{1}{2p^2(1 - \cos \theta) + \mu^2} + \frac{1}{2p^2(1 + \cos \theta) + \mu^2} \right).$$